

A New Class of Well-Balanced Finite Volume schemes for Conservation laws with source terms

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Joint Work with:

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- ▶ Nils Henrik Risebro (CMA, Oslo).

The Problem

Numerical Difficulties

Existing Well-Balanced Schemes

New Well-Balanced Schemes

Numerical Experiments

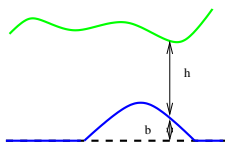
Summary and Future Work

Basic Equations

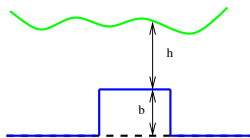
$$U_t + (f(U))_x + (g(U))_y + (h(U))_z = S(x, U)$$

- ▶ System of Conservation laws in multi-D.
- ▶ Together with source terms.
- ▶ Source can be spatially dependent (maybe singular).
- ▶ Also termed **Balance** laws.

Flow on a non-trivial topography



Non-Trivial Smooth Bottom Topography



Discontinuous Bottom Topography

An Example

- Shallow water equations with Non-trivial Bottom Topography.

$$\begin{aligned} h_t + (hu)_x + (hv)_y &= 0 \\ (hu)_t + (hu^2 + \frac{1}{2}gh^2)_x + (huv)_y &= -ghb_x \\ (hv)_t + (huv)_x + (hv^2 + \frac{1}{2}gh^2)_y &= -ghb_y \end{aligned}$$

- ▶ h is height of the free surface.
- ▶ (u, v) is the velocity vector.
- ▶ g - gravity constant.
- ▶ b Topography function (can be **discontinuous**).

The Model Equation

- **Single conservation law in 1-d.**

$$u_t + f(u)_x = A(x, u)$$

- ▶ Unknown u , flux f and source A .
- ▶ Source can even be singular. (A can be a measure)

Special Cases

- *Autonomous source*

$$u_t + f(u)_x = g(u)$$

- *Scalar “Shallow Water” equations*

$$u_t + (f(u))_x = z'(x)b(u)$$

- ▶ z is the topography function (possibly discontinuous)

- *Singular Sources*

$$u_t + (f(u))_x = z'(x)$$

- z Heaviside function \Rightarrow RHS is a measure.

Weak Solutions

- Well Defined when $A(x, u) \in L^\infty$.
- $u \in L^\infty(\mathbb{R} \times \mathbb{R}_+) \cap L^1_{loc}$ is a weak solution if for all test functions φ ,

$$\int_{\mathbb{R}^+} \int_{\mathbb{R}} u \varphi_t + f(u) \varphi_x + A(x, u) \varphi \, dx dt + \int_{\mathbb{R}} u(x, 0) \varphi(x, 0) \, dx = 0 \quad (1)$$

- Special attention when $A \notin L^\infty$.
- Make sense of the non-conservative product

$$z'(x)b(u)$$

Entropy Solutions

- Well Defined when $A(x, u) \in L^\infty$.
- $u \in L^\infty(\mathbb{R} \times \mathbb{R}_+) \cap L^1_{loc}$ is a entropy solution if for all test functions $\varphi \geq 0$,

$$\int_{\mathbb{R}^+} \int_{\mathbb{R}} S(u) \varphi_t + Q(u) \varphi_x + S'(u) A(x, u) \varphi \, dx dt + \int_{\mathbb{R}} u(x, 0) \varphi(x, 0) \, dx \geq 0$$

- For any entropy-entropy flux pair (S, Q) .
- Entropy solutions exist and are unique when $A \in L^\infty$.
- No general theory in the singular case except when $A(x, u) = z'(x)$

A Naive Numerical Scheme

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- ▶ Explicit Euler in Time.
- ▶ Godunov type numerical fluxes for the flux.
- ▶ Central differences for the source.

A Numerical Experiment

$$u_t + f(u)_x = z'(x)b(u)$$

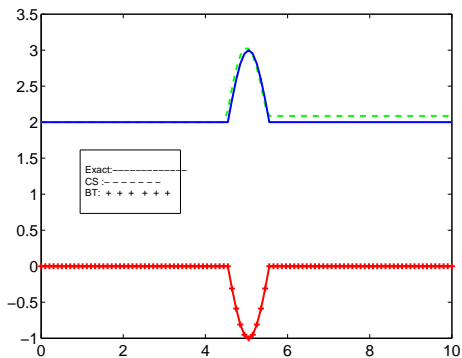
- With

$$\begin{aligned}
 f(u) &= \frac{1}{2}u^2 & b(u) &= u \\
 -z(x) &= \begin{cases} \cos(\pi x) & \text{if } 4.5 < x < 5.5 \\ 0 & \text{Otherwise} \end{cases} \\
 u(t, 0) &= 2 & u(0, x) &= 0
 \end{aligned}$$

- Explicit steady state is given by

$$\bar{u}(x) = 2 + z(x)$$

At the steady state



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$$f(u)_x \approx A(x, u)$$

- ▶ Numerical schemes have to preserve Flux-Source balance.
- ▶ Centered Source/Operator splitting doesn't respect it.
- ▶ Search for better schemes

- ▶ Greenberg, Leroux.
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- ▶ Gosse, Leroux.
- ▶ Botchorischvili, Perthame and Vasseur. (BPV)
- ▶ Bermudez, Vasquez
- ▶ Perthame, Bouchut, Bristeau, Klien, Audusse.
- ▶ Russo, Noelle, Kurganov, Levy and many others.

WBS (condensed)

- Consider Scalar Shallow water equations,

$$u_t + f(u)_x = z'(x)b(u)$$

- The steady state is formally,

$$\begin{aligned} f(u)_x &= z'(x)b(u) \\ \Rightarrow f'(u)u_x &= z'(x)b(u) \\ \Rightarrow \frac{f'(u)}{b(u)}u_x &= z'(x) \\ \Rightarrow D(u)_x &= z'(x) \\ D(u) &= \int^u \frac{f'(s)}{b(s)} ds \end{aligned}$$

- Steady State evaluated from

$$D - z = \text{Constant}$$

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- ▶ At n th time step let v_j^n be the cell-averages and z_j be averages of the topography, then define “local” steady states solving

$$\begin{aligned} D(v_j^n -) - z_j &= D(v_{j-1}^n) - z_{j-1} \\ D(v_j^n +) - z_j &= D(v_{j+1}^n) - z_{j+1} \end{aligned}$$

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- ▶ Use the local steady states to define a Godonov type scheme with update
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$$v_j^{n+1} = v_j^n - \frac{\Delta t}{\Delta x} (F(v_j^n, v_j^n+) - F(v_j^n-, v_j^n))$$

- with F being Standard (Godunov, Enquist-Osher) flux corresponding to f

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- ▶ Shown to Converge to entropy solutions (via Kinetic formulation).
- ▶ Basis for WBS for systems.

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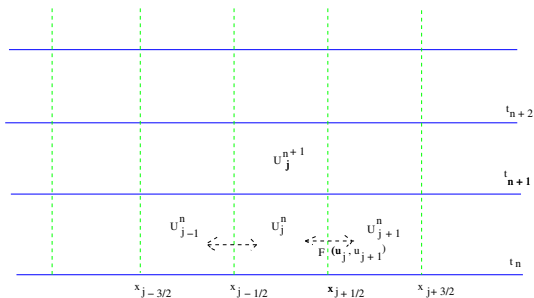
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- ▶ Possible loss of accuracy away from steady states.
- ▶ Subtle deficiencies (see sequel)

Along the lines of

- ▶ Greenberg, Leroux, Baraille and Noussair. (Singular Sources)
- ▶ Noussair.
- ▶ LeVeque.
- ▶ Bale, LeVeque, Mitran, Rossmanith (Flux - Differencing)
- ▶ Adimurthi, Gowda, Mishra (Singular Sources)

1-D Finite Volume Grid



The Scheme: Design

- ▶ At time level n , let u_j^n be the cell averages,
 - Step 1: Freeze the source at t^n and define the piecewise constant

$$u^n(x) = \sum_j u_j^n \mathbf{1}_{\{I_j\}}(x)$$

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- ▶ Primitive Reconstruction: Define the function

$$\tilde{B}^n(x) = \int^x A(y, u^n(y)) dy$$

- We obtain the following discontinuous flux problem,

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- ▶ Local Discontinuous flux problems: By sampling define

$$B^n(x) = \sum_j \tilde{B}^n(x_j) \mathbf{1}_{\{I_j\}}(x)$$

- We obtain the following discontinuous flux problem,

$$u_t + (f(u) - B^n(x))_x = 0, u(x, t^n) = u^n(x)$$

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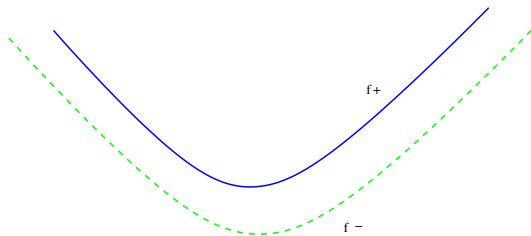
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- ▶ Local Riemann problems at each interface

$$\begin{array}{ll} u_t + (f(u) - B_j^n)_x = 0 & u(x, 0) = u_j^n \quad x < x_{j+1/2} \\ u_t + (f(u) - B_{j+1}^n)_x = 0 & u(x, 0) = u_{j+1}^n \quad x > x_{j+1/2} \end{array}$$

Shape of Adjacent fluxes



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$$u_j^{n+1} = u_j^n - \frac{\Delta t}{\Delta x} (F_{j+1/2}^n - F_{j-1/2}^n)$$

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- with $F_{j+1/2}$ being the corresponding Godunov flux.
- ▶ Explicit formulas are available in most cases e.g (f convex) then

$$F_{j+1/2} = \max(f(\max(u_j, \theta)) - B_j^n, f(\min(u_{j+1}, \theta)) - B_{j+1}^n)$$

Scheme: Properties

- ▶ Discrete steady states of the scheme

$$f(u_{j+1}^n) - f(u_j^n) = B_{j+1}^n - B_j^n$$

- Reflects Flux-Source balance.

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- Reflects Flux-Source balance.
- ▶ Rankine-Hugoniot Conditions + Jump entropy conditions \Rightarrow *Entropic* Discrete steady states are preserved .
- ▶ Flexibility in the averaging steps to obtain equivalent discrete steady states.

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Approximations (**Compensated Compactness**).
- ▶ Jump entropy conditions + Structure of the scheme \Rightarrow
Convergence to entropy solutions if $A \in L^\infty$

Experiment 1 (Continuous Bottom)

$$u_t + f(u)_x = z'(x)b(u)$$

- With

$$f(u) = \frac{1}{2}u^2$$

$$-z(x) = \begin{cases} \cos(\pi x) & \text{if } 4.5 < x < 5.5 \\ 0 & \text{if Otherwise} \end{cases}$$

$$u(t, 0) = 2$$

$$b(u) = u$$

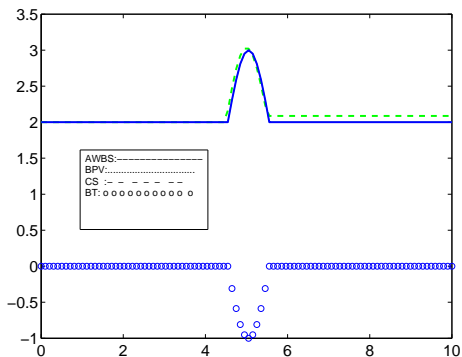
$$u(0, x) = 0$$

- Explicit steady state is given by

$$\bar{u}(x) = 2 + z(x)$$

- Comparison of Central Sources (CS), Existing Well-Balanced Scheme (BPV) and New Well-Balanced Scheme (AWBS)

At Steady State: $\Delta x = 0.1$



Errors at Steady State

	L^∞	L^1
CS	0.1652	0.4824
AWBS	4.37×10^{-14}	2.22×10^{-13}
BPV	8.45×10^{-14}	2.26×10^{-13}

Transients: $\Delta x = 0.1$

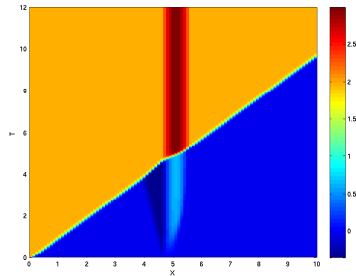
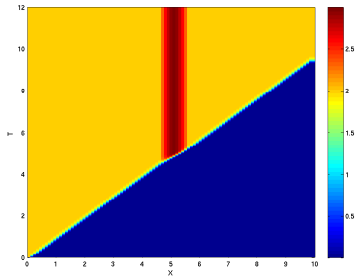
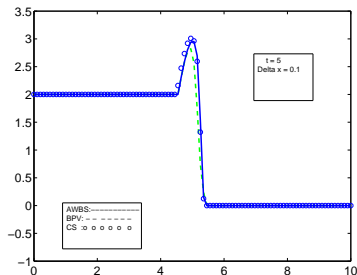
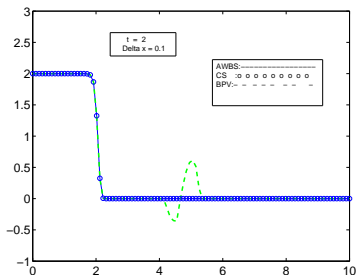
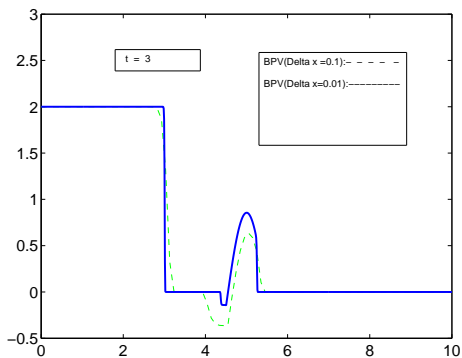


Figure: Left: **AWBS**, Right: **BPV**

Transient snapshots



BPV at high resolution



Experiment 2 (Discontinuous Bottom)

$$u_t + f(u)_x = z'(x)b(u)$$

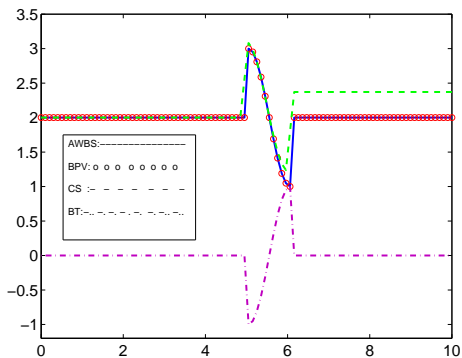
- With

$$\begin{aligned}
 f(u) &= \frac{1}{2}u^2 & b(u) &= u \\
 -z(x) &= \begin{cases} \cos(\pi x) & \text{if } 5 < x < 6 \\ 0 & \text{if Otherwise} \end{cases} \\
 u(t, 0) &= 2 & u(0, x) &= 0
 \end{aligned}$$

- Explicit steady state is given by

$$\bar{u}(x) = 2 + z(x)$$

At Steady State: $\Delta x = 0.1$

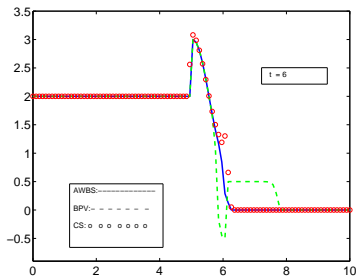
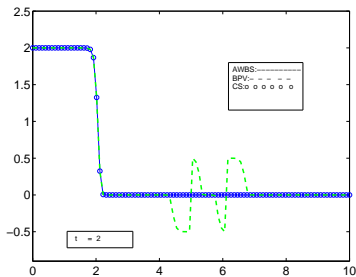


Errors at Steady State

	L^∞	L^1
CS1	0.8027	1.6449
AWBS	1.87×10^{-12}	8.12×10^{-13}
BPV	2.53×10^{-9}	6.34×10^{-10}

Table: Errors at the steady state for the three schemes with $\Delta x = 0.1$ in Experiment 2

A Transient snapshot



Experiment 3 (Another Discontinuous Bottom)

$$u_t + f(u)_x = z'(x)b(u)$$

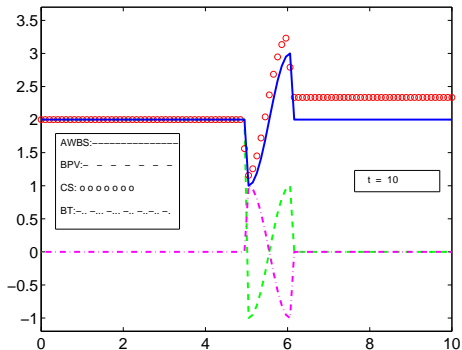
- With

$$\begin{aligned}
 f(u) &= \frac{1}{2}u^2 & b(u) &= u \\
 -z(x) &= \begin{cases} -\cos(\pi x) & \text{if } 5 < x < 6 \\ 0 & \text{if Otherwise} \end{cases} \\
 u(t, 0) &= 2 & u(0, x) &= 0
 \end{aligned}$$

- Explicit steady state is given by

$$\bar{u}(x) = 2 + z(x)$$

At Steady State: $\Delta x = 0.1$



Transients

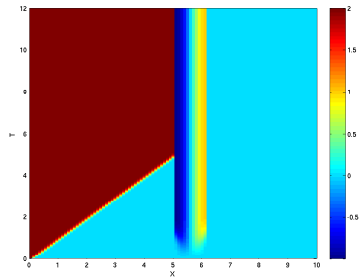
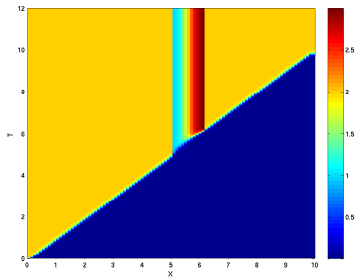


Figure: Left: **AWBS**,

Right: **BPV**

Experiment 4 (Non-Monotone D)

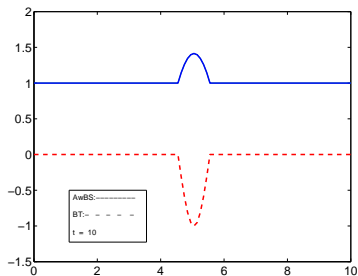
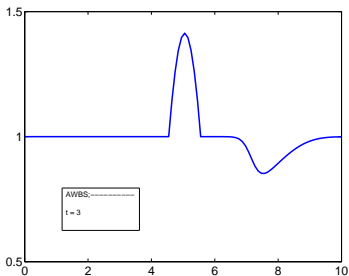
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$$\begin{aligned}
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 u(t, 0) &= 2 & u(0, x) &= 0
 \end{aligned}$$

- Difficult to define **BPV** as $D = u^2$ is not monotone. No problems with **AWBS**

Numerical Results at : $\Delta x = 0.1$



Experiment 5 (Source not in Product form)

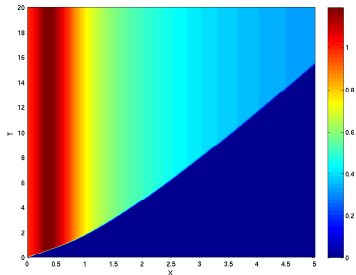
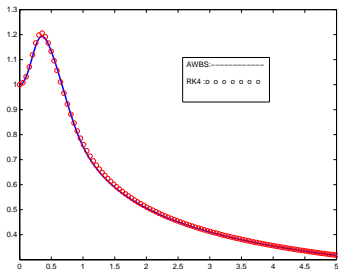
$$u_t + f(u)_x = A(x, u)$$

- With

$$\begin{aligned} f(u) &= \frac{1}{2}u^2 & A(x, u) &= \sin(2\pi x u^2) \\ u(t, 0) &= 1 & u(0, x) &= 0 \end{aligned}$$

- Unclear how to define **BPV** in this case (other than using ODE solvers at each mesh point) whereas **AWBS** is well-defined

Numerical results with : $\Delta x = 0.1$



Subtle problems with existing well-balanced schemes

- ▶ Incorrect Shock speeds and strengths due to non-linear transformations.

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- ▶ Incorrect Shock speeds and strengths due to non-linear transformations.
- ▶ Problems at resonance $u = 0$

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 - ▶ Very Simple to implement (Explicit formulas, No extra equations)
 - ▶ Robust and proved to be convergent.
 - ▶ Very General: Work with different type of fluxes and sources.

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